## Advanced Statistical Physics II – Problem Sheet 12

## Problem 1 – Many-Particle Non-Equilibrium System

In the lecture, we discussed a class of non-equilibrium systems whose dynamics can be modelled by the following multi-dimensional Langevin equation

$$\dot{z}_k(t) = -A_{km} z_m(t) + \Phi_{km} F_m(t), \quad \langle F_m(t) F_n(t') \rangle = 2\delta_{mn} \delta(t - t') \quad . \tag{1}$$

Note, that we make use of the Einstein summation convention for double indices. The corresponding Fokker-Planck equation reads

$$\dot{P}(\vec{z},t) = \left[\nabla_k A_{km} z_m + \nabla_k \nabla_m C_{km}\right] P(\vec{z},t), \quad C_{ij} := \Phi_{ik} \Phi_{jk} \quad .$$
<sup>(2)</sup>

a) (5P) Use a Gaussian ansatz  $P_0(\vec{z}) = \mathcal{N}^{-1} \exp(-z_i E_{ij}^{-1} z_j/2)$ , for the stationary solution. Here  $E_{ij} = \langle z_i z_j \rangle$  denotes the entries of the symmetric covariance matrix. Derive the Lyapunov equation discussed in the lecture:

$$A_{ik}E_{kj} + A_{jk}E_{ki} = 2C_{ij} \tag{3}$$

b) (5P) As a concrete example, we will discuss a simple system of an overdamped particle in harmonic confinement of strength *M*. Furthermore, it is harmonically coupled to a second particle:

$$H(x,y) = \frac{M}{2} x^{2} + \frac{K}{2} (x-y)^{2}$$
(4)

Both particles are subject to friction  $\gamma_{x,y}$  and noise of strength  $b_{x,y}$ . This leads to

$$z(t) = \begin{pmatrix} x(t) \\ y(t) \end{pmatrix}, \quad A = \begin{pmatrix} (K+M)/\gamma_x & -K/\gamma_x \\ -K/\gamma_y & K/\gamma_y \end{pmatrix}, \quad \Phi = \begin{pmatrix} b_x/\gamma_x & 0 \\ 0 & b_y/\gamma_y \end{pmatrix},$$
(5)

In equilibrium, friction and noise strength are related via  $b_x^2/\gamma_x = b_y^2/\gamma_y = k_B T$ . Formally, we can understand departure from equilibrium by introducing different temperatures for each particle  $k_B T_x = b_x^2/\gamma_x$  and  $k_B T_y = b_y^2/\gamma_y$ . Find the entries of the covariance matrix

$$E = \begin{pmatrix} \langle x^2 \rangle & \langle xy \rangle \\ \langle xy \rangle & \langle y^2 \rangle \end{pmatrix}$$
(6)

by solving the Lyapunov equation (3). Express the resulting covariances in terms of the dimension-less parameter

$$\alpha := \frac{T_y - T_x}{T_x} = \frac{\gamma_x}{b_x^2} \cdot \frac{b_y^2}{\gamma_y} - 1 \tag{7}$$

which quantifies departure from equilibrium, and in terms of  $T_x$ . Check that in equilibrium ( $\alpha = 0$ ) the equipartition theorem for the variables x and x - y is obeyed.

c) (3P) Consider a hypothetical experiment in which only the position *x* can be observed. Thus the position *y* of the other particle is a hidden degree of freedom. An example would be a colloid in a laser trap whose position is tracked. In equilibrium, the position of the trapped particle obeys a Boltzmann distribution

$$P_{\rm eq}(x) \propto e^{-Mx^2/2k_BT} \tag{8}$$

i.e. the variance is given by  $\langle x^2 \rangle = k_B T/M$ . How does the variance/width of the observed distribution change for  $T_x > T_y$  and  $T_x < T_y$ ?

## Problem 2 – Run and Tumble Particle

Consider the following dynamics of a free, overdamped particle in *d* dimensions:

$$\dot{\vec{x}}(t) = \vec{u}(t) + \gamma^{-1}\vec{F}_R(t), \qquad \langle \vec{F}_R(t)\vec{F}_R(t') \rangle = 2dk_B T \gamma \delta(t-t')$$
(9)

Here,  $\gamma$  denotes the friction coefficient and  $\vec{F}_R(t)$  is the random force, which accounts for thermal fluctuations. The particle propels itself forward at constant velocity  $|\vec{u}(t)| = v_0$ . (For  $v_0 = 0$  (and thus  $\vec{u}(t) = \vec{0}$ ), this corresponds to standard Brownian motion.) The particle goes in the same direction for an average time  $\tau$ and then chooses a new direction completely at random - independent of the previous orientation and thermal noise. Assuming the time for going in one direction is exponentially distributed, the autocorrelation of  $\vec{u}(t)$  is given by

$$\langle \vec{u}(t)\vec{u}(t')\rangle = v_0^2 e^{-|t-t'|/\tau}$$
(10)

This model has been proposed for the dynamics of bacterial motility. Fig. 1 shows a simulated trajectory of such a run and tumble particle.



Figure 1: Trajectory of a run and tumble particle in two dimensions.

- a) (4P) Calculate the mean-square displacement  $\Delta x^2(t) = \langle (x(t) x(0))^2 \rangle$ . Note that you need to average over both  $\vec{F}_R(t)$  and  $\vec{u}(t)$ . *Hint*: It is helpful to first calculate the velocity autocorrelation function  $\langle \vec{x}(t)\dot{\vec{x}}(t') \rangle$  and then obtain the mean-square displacement via integration.
- b) (3P) For both short and long times *t*, the dynamics is diffusive. Calculate the diffusion constants for short and long times:

$$D_{\text{short}} = \lim_{t \to 0} \frac{\Delta x^2(t)}{2t}, \qquad D = \lim_{t \to \infty} \frac{\Delta x^2(t)}{2t}$$
(11)

Interpret the result. How does the active self propulsion affect the diffusion constant?