Second Chance Exam Solutions: Advanced Statistical Physics Part II: Problems (75P)

1 Three-Spin Interaction (20P)

Consider a one dimensional system of N spins with the following Hamiltonian:

$$H = -J \sum_{i=1}^{N-2} S_i S_{i+1} S_{i+2}$$

with $S_i = \pm 1, i = 1...N$ being the spin states and J = const being an interaction parameter. Assume open boundary conditions (**not** periodic!).

Calculate the canonical partition function and the Helmholtz free energy F. What is the thermodynamic limit of F?

Hint: The transfer-matrix method is not necessary for this problem.

The definition of the cosinus hyperbolicus may be helpful:

$$\cosh(x) = \frac{e^x + e^{-x}}{2}$$

$$Z_{N} = \sum_{S_{1},...,S_{N}} e^{\beta J \sum_{i=1}^{N-2} S_{i} S_{i+1} S_{i+2}}$$

= $\sum_{S_{1},...,S_{N-1}} e^{\beta J \sum_{i=1}^{N-3} S_{i} S_{i+1} S_{i+2}} \sum_{S_{N}} e^{\beta J S_{N-2} S_{N-1} S_{N}}$
= $\sum_{S_{1},...,S_{N-1}} e^{\beta J \sum_{i=1}^{N-3} S_{i} S_{i+1} S_{i+2}} 2 \cosh(\beta J)$
= $\sum_{S_{1},S_{2},S_{3}} e^{\beta J S_{1} S_{2} S_{3}} (2 \cosh(\beta J))^{N-3}$
= $4 (2 \cosh(\beta J))^{N-2}$

Therefore the Helmholtz free energy is:

$$F = -\frac{1}{\beta} \ln Z_N = -\frac{1}{\beta} \Big(\ln 4 + (N-2) \ln(2 \cosh(\beta J)) \Big)$$

For the thermodynamic limit we get:

$$\lim_{N\to\infty} F = -\frac{N}{\beta} \, \ln(2 \, \cosh(\beta \, J))$$

2 Liquid-Gas Phase-Transition (25P)

We want to consider a substance with the enthalpy of the liquid phase

$$H_l(p, S) = 2 (a \, p \, S \, N)^{1/2}$$

and the enthalpy of the gas phase

$$H_g(p,S) = 3\left(\frac{pS}{2}\right)^{2/3} (bN)^{1/3},$$

where a > 0 and b > 0 are constants, S is the total entropy of the system, p is the pressure and N is the total number of particles.

a) Calculate the liquid-gas coexistence temperature T_g as a function of pressure. (16P)

We have to perform a Legrende-Transformation from H(S,p) to G(T,p) = H - TS with $T = \partial H / \partial S$: liquid phase:

$$T = \partial H_l / \partial S = \sqrt{\frac{aNp}{S}}$$
(1)

$$\Rightarrow S = \frac{\alpha \cdot p}{T^2} \tag{2}$$

$$G_{l}(T,p) = 2aN\frac{1}{T} - aN\frac{1}{T} = aN\frac{1}{T}$$

$$g_{l}(T,p) = G_{l}/N = a\frac{p}{T}$$
(3)
(4)

gas phase:

$$T = \partial H_g / \partial S = 2 \left(\frac{p}{2}\right)^{2/3} \left(\frac{bN}{S}\right)^{1/3}$$
(5)

$$\Rightarrow S = 2bN\frac{p}{T^3} \tag{6}$$

$$G_g(T,p) = 3bN\frac{p^2}{T^2} - 2bN\frac{p^2}{T^2} = bN\frac{p^2}{T^2}$$
(7)

$$g_g(T,p) = G_g/N = b\frac{p^2}{T^2}$$
(8)

At the phase coexistence line we have:

$$g_l(T_v, p_v) = g_g(T_v, p_v)$$

$$(9)$$

$$n_v = n^2$$

$$\Rightarrow a \frac{p_v}{T_v} = b \frac{p_v^2}{T_v^2} \tag{10}$$

$$\Rightarrow T_v = p_v \frac{\dot{b}}{a} \tag{11}$$

b) Calculate the densities of the liquid and the gas phase at the phase transition line. (6P) liquid phase:

$$V_l = \partial G_l / \partial p = \frac{aN}{T_v}$$
(12)
$$\Rightarrow c_v = N/V = \frac{T_v}{T_v}$$
(13)

$$\Rightarrow \rho_l = N/V_l = \frac{T_v}{a} \tag{13}$$

gas phase:

$$V_g = \partial G_g / \partial p = \frac{2bNp_v}{T_v^2} \stackrel{(11)}{=} \frac{2aN}{T}$$
(14)

$$\Rightarrow \rho_g = N/V_g = \frac{T_v}{2a} \tag{15}$$

c) Calculate the entropy change per volume $\Delta S/\Delta V$ at the phase transition line. (3P) We can use Clausius-Clapeyron:

$$\frac{\Delta S}{\Delta V} = \frac{\Delta S/N}{\Delta V/N} = \Delta s/\Delta v = \frac{\mathrm{d}p_v}{\mathrm{d}T_v} \stackrel{(11)}{=} \frac{a}{b} \tag{16}$$

3a Wien's Law in n Dimensions (15P)

 $u(\omega)$ is the energy density of black body radiation at angular frequency ω per volume and angular frequency, i.e. the spectral density of the internal energy density U/V.

Determine $u(\omega)$ in n dimensions in the low temperature limit.

The volume of an *n*-dimensional sphere is $C_n \mathbb{R}^n$, where \mathbb{R} is the radius. You do not have to determine C_n .

$$U = 2\frac{V}{h^n} \int d^n p \, \frac{e^{-\beta pc}}{1 - e^{-\beta pc}} pc \tag{17}$$

$$= 2\frac{V}{h^n}C_n \int \mathrm{d}p \,\frac{e^{-\beta pc}}{1 - e^{-\beta pc}} n p^{n-1} pc \tag{18}$$

$$= 2n \frac{cV}{h^n} C_n \int \mathrm{d}p \, \frac{e^{-\beta pc}}{1 - e^{-\beta pc}} p^n \tag{19}$$

$$pc = \hbar\omega \Rightarrow \mathrm{d}p = \frac{\hbar}{c}\mathrm{d}\omega$$
 (20)

$$\Rightarrow U = 2n \frac{cV}{h^n} C_n \int d\omega \, \frac{\hbar}{c} \frac{e^{-\beta\hbar\omega}}{1 - e^{-\beta\hbar\omega}} \left(\frac{\hbar\omega}{c}\right)^n \tag{21}$$

$$= 2n \frac{V\hbar}{(2\pi c)^n} C_n \int d\omega \, \frac{e^{-\beta\hbar\omega}}{1 - e^{-\beta\hbar\omega}} \omega^n \tag{22}$$

For $\beta \hbar \omega \gg 1$ we can approximate $1 - e^{-\beta \hbar \omega} \approx 1$:

$$\Rightarrow U = 2n \frac{V\hbar}{(2\pi c)^n} C_n \int d\omega \, e^{-\beta\hbar\omega} \omega^n$$
(23)

$$\frac{U}{V} = \int d\omega \, u(\omega) \tag{24}$$

$$\Rightarrow u(\omega) = 2n \frac{\hbar}{(2\pi c)^n} C_n e^{-\beta\hbar\omega} \omega^n$$
(25)

3b Cosmic Background Radiation (15P)

The universe is filled with a photon gas that corresponds to black body ratiation of temperature $T_{present} = 3 \text{ K}$. In a simple view, this radiation arose from the isentropic expansion of a much hotter photon cloud, which was produced during the big bang.

If the volume of the universe, and thus the volume of the photon gas, increases isentropically by a factor of two starting from the present state, what will be the final temperature of the photon gas?

Hint: The Stefan-Boltzmann law can be useful.

Stefan-Boltzmann law: $U\propto VT^4$

The total differential is

$$dU = TdS + pdV$$

$$(26)$$

$$(\partial U) \quad (\partial U) \quad (\partial T)$$

$$\Rightarrow T = \left(\frac{\partial C}{\partial S}\right)_{V} = \left(\frac{\partial C}{\partial T}\right)_{V} \left(\frac{\partial I}{\partial S}\right)_{V}$$
(27)
(*∂S*) = 1 (*∂U*) (27)

$$\Rightarrow \left(\frac{\partial S}{\partial T}\right)_{V} = \frac{1}{T} \left(\frac{\partial C}{\partial T}\right)_{V} \tag{28}$$

$$\Rightarrow S \propto VT^3 \tag{29}$$

is entropic: $\mathrm{dS}=0$

$$\Rightarrow V_1 T_1^3 = V_2 T_2^3$$

$$\Rightarrow T_2 = T_1 (V_1 / V_2)^{1/3} = T_1 2^{-1/3}$$
(30)
(31)